

A Gentle Introduction to Magnetic Reconnection

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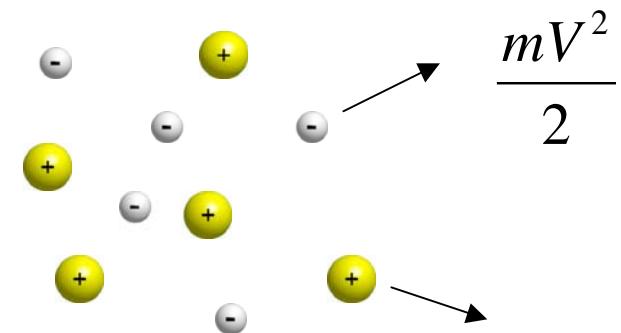
What is Magnetic Reconnection?

Magnetic Energy



Kinetic Energy

$$\begin{array}{c} \rightarrow \\ \rightarrow \\ \rightarrow \\ \rightarrow \end{array} \frac{B^2}{8\pi}$$



Many applications

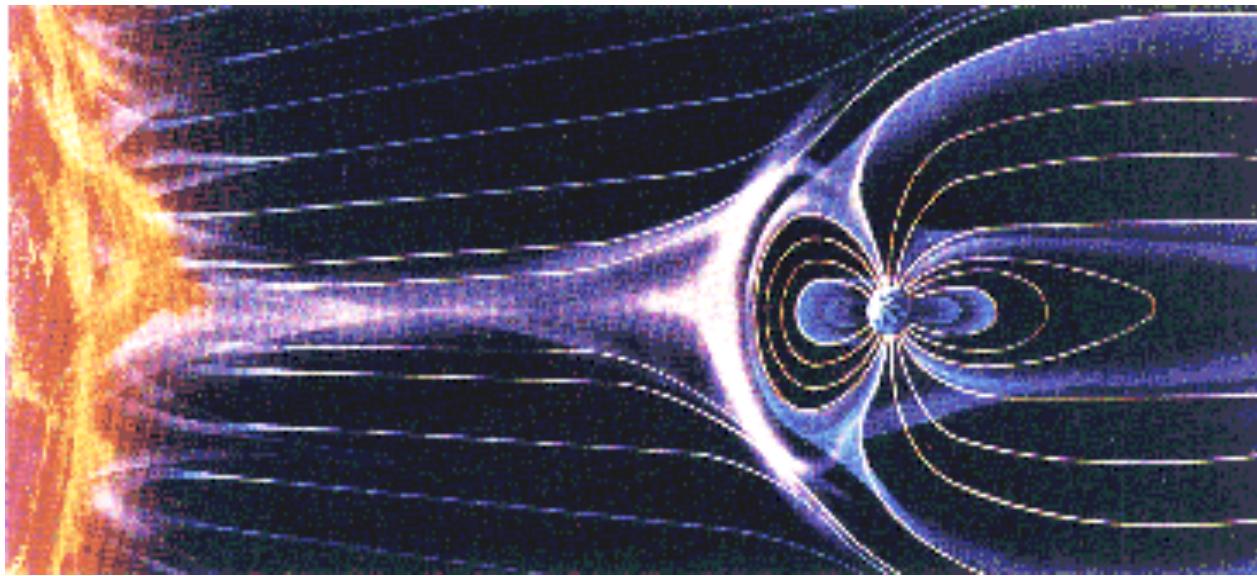
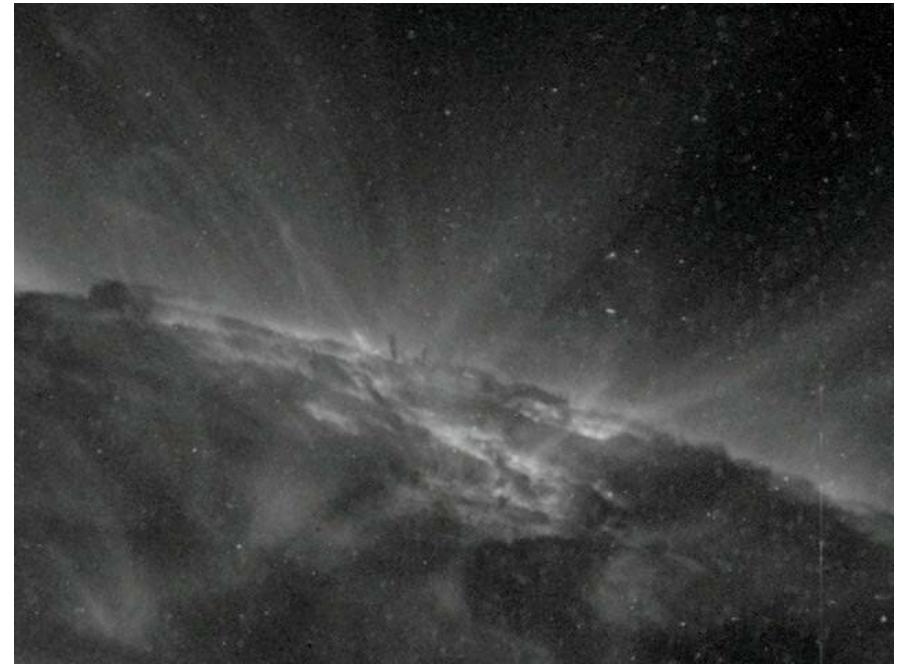


Space Plasmas

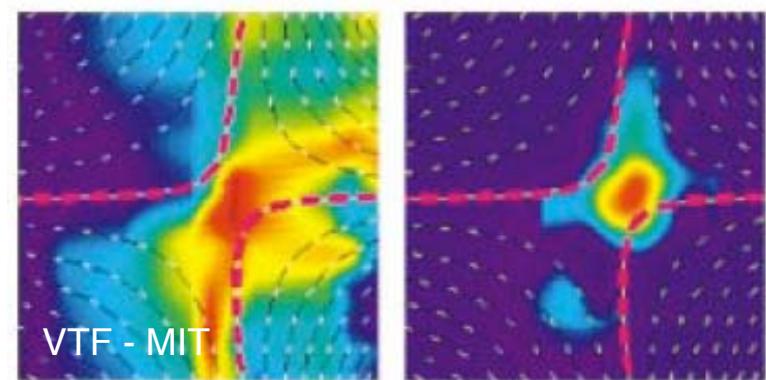
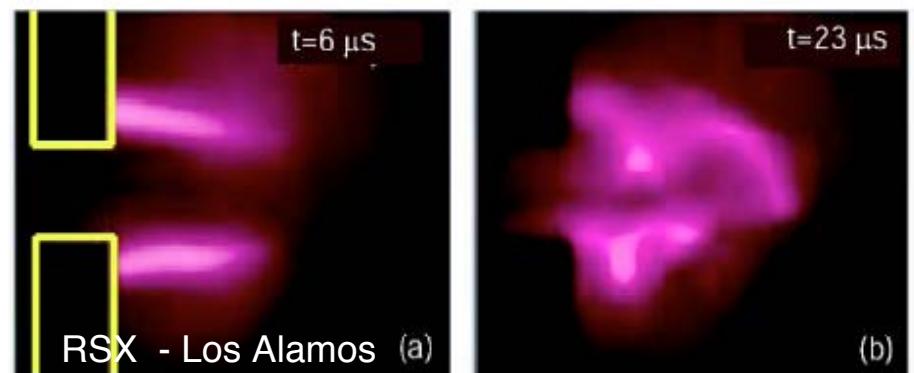
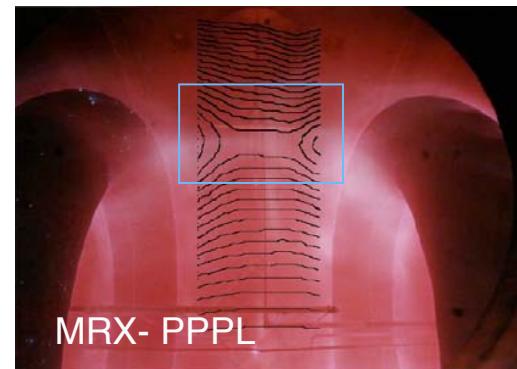
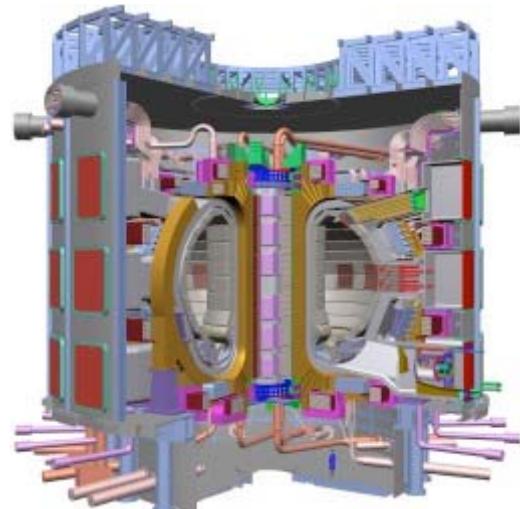
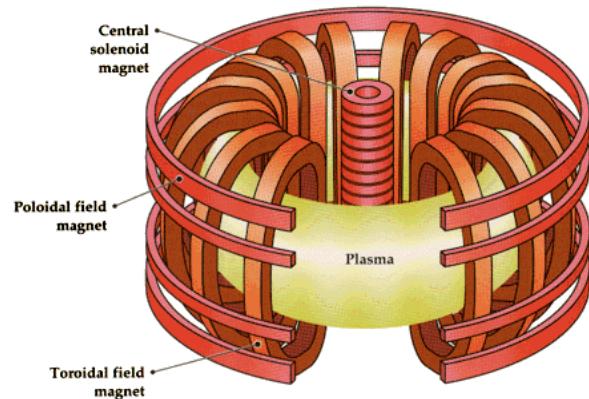
Laboratory Experiments

Astrophysical Plasmas

Reconnection in Space Plasmas

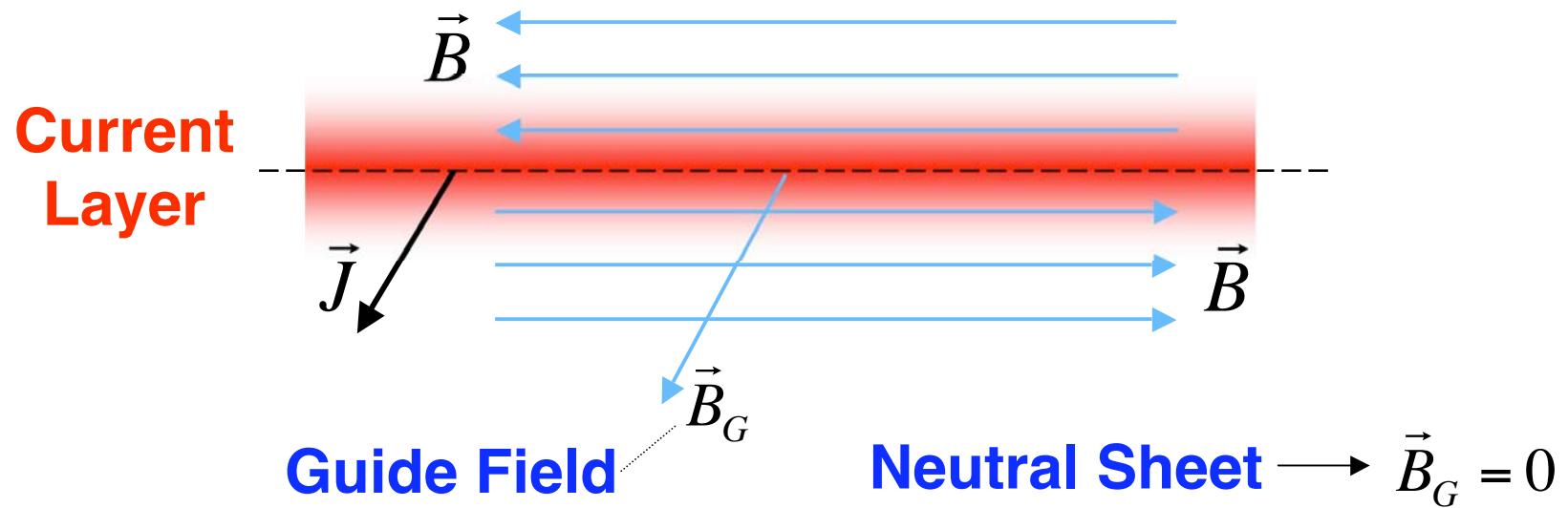


Reconnection in Laboratory Plasmas

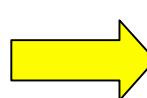


What is a Current Sheet?

Current layer + corresponding field reversal



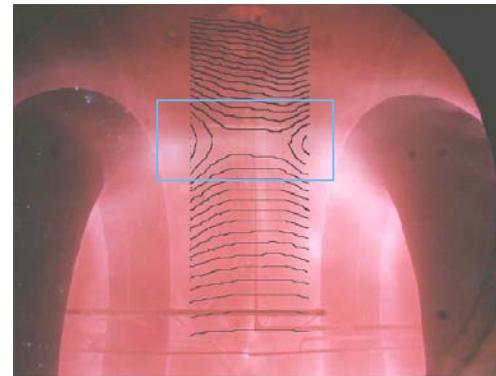
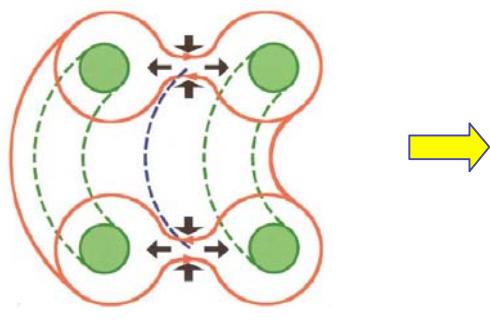
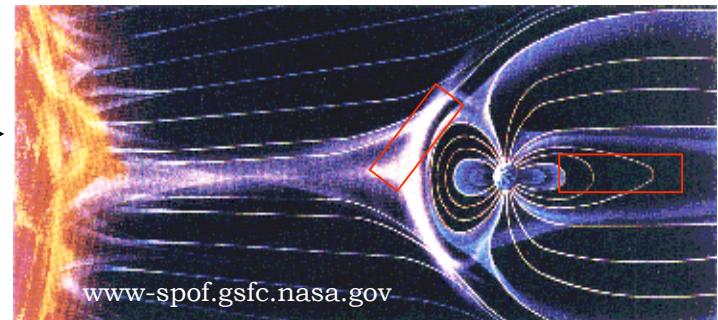
SciSearch Database
“current + sheet”



$\sim 10,000$ papers
since 1960

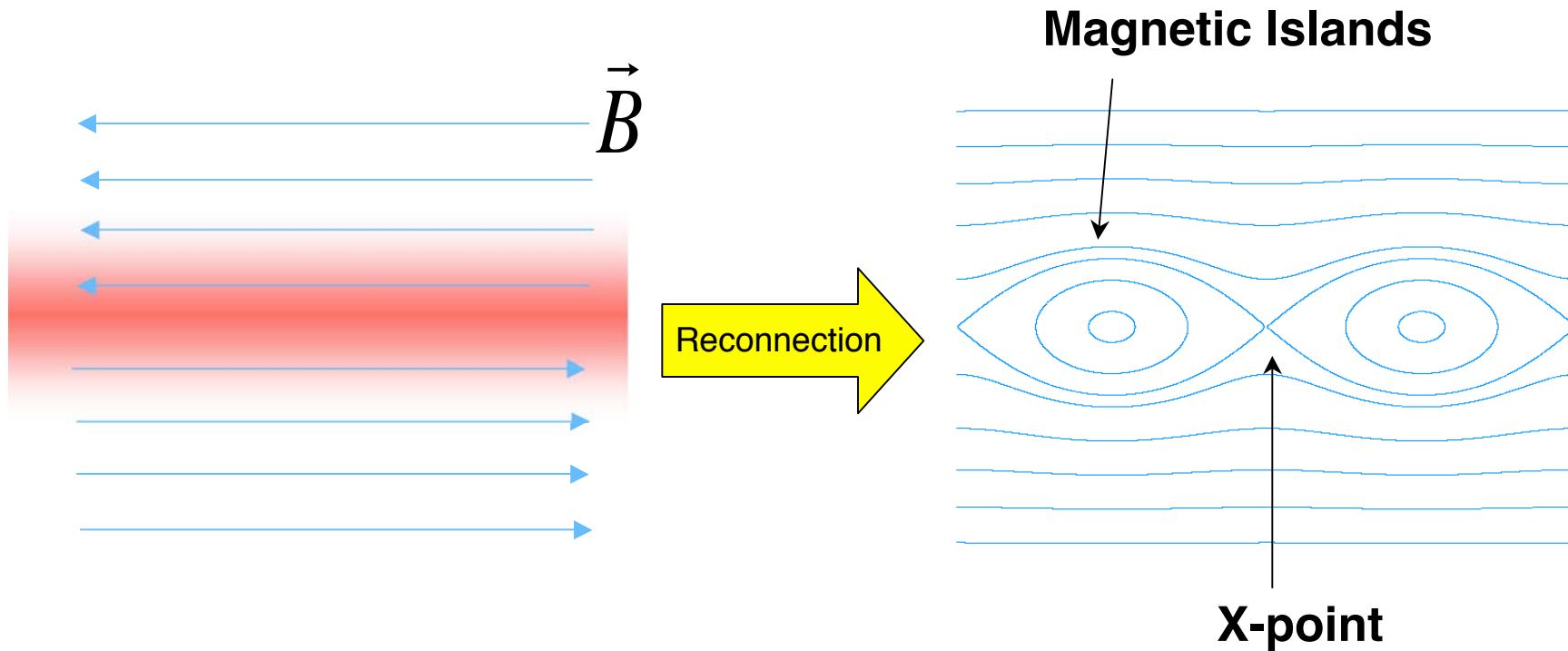
Examples of Current Sheets

- Magnetotail
- Magnetopause
- Heliospheric current sheet
- Plasma tail of comets
- Solar flares & prominences
- Simple geometry to study magnetic reconnection
- Laboratory plasmas – MRX experiment at PPPL

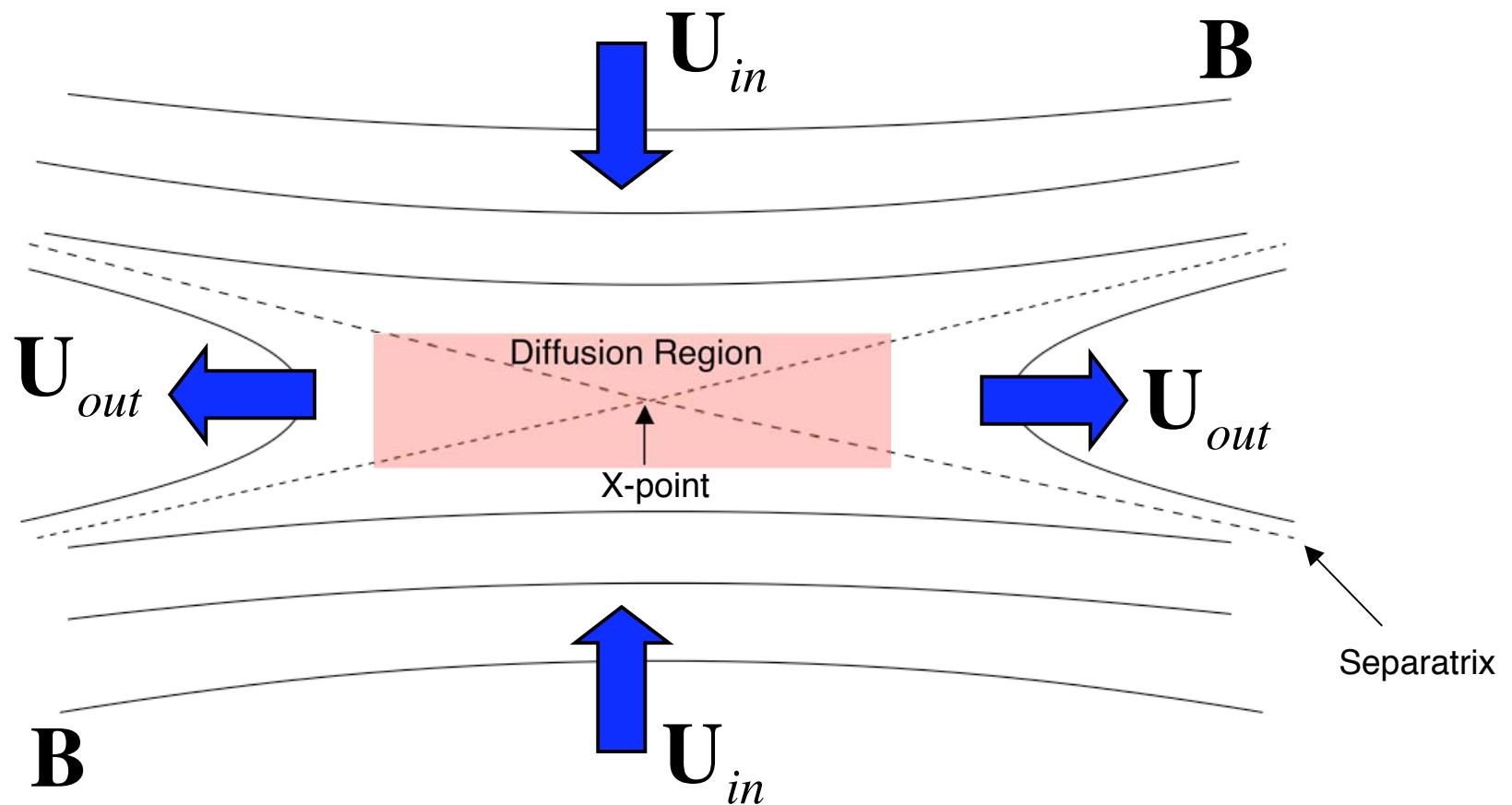


Figures courtesy of Hantao Ji (PPPL)

Topological Consequences



More Terminology



Collisional vs Collisionless Reconnection

$$\nu \propto \frac{n}{T^{3/2}}$$

**Magnetotail
Parameters**

$$T_e \sim 1 \text{ keV} \quad T_i \sim 6 \text{ keV} \longrightarrow \nu_e \sim 10^{-9} \text{ sec}^{-1} \quad \Omega_{ci} \sim 1 \text{ sec}^{-1}$$
$$n \sim 1 \text{ cm}^{-3} \quad B \sim 20 \text{ nT} \quad \frac{V_{the}}{\nu_e L} \sim 10^{10} \quad \frac{\rho_i}{L} \sim 1$$
$$L \sim 1000 \text{ km}$$

Questions for Collisionless Regime:

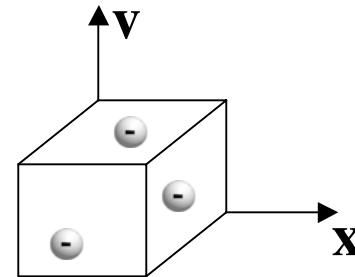
1. How does reconnection proceed so rapidly in collisionless regimes?

2. How does it get started in the first place?  **Onset problem**

What equations describe a plasma?

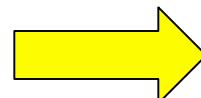
$$\frac{\partial f_s}{\partial t} + \mathbf{v} \cdot \frac{\partial f_s}{\partial \mathbf{x}} + \frac{q_s}{m_s} \left(\mathbf{E} + \frac{\mathbf{v} \times \mathbf{B}}{c} \right) \cdot \frac{\partial f_s}{\partial \mathbf{v}} = 0 \quad \xleftarrow{\epsilon = 0} \text{Vlasov}$$

$f_s(\mathbf{x}, \mathbf{v}, t) \rightarrow \frac{\text{Number of particles}}{\text{Unit volume of phase space}}$



Small Parameter $\rightarrow \epsilon = \frac{1}{n\lambda_D^3} \sim 10^{-6} \rightarrow 10^{-12}$

$$\rho = \sum_s q_s \int f_s d\mathbf{v}$$



$$\mathbf{J} = \sum_s q_s \int \mathbf{v} f_s d\mathbf{v}$$

Maxwell's Equations

$$\nabla \cdot \mathbf{B} = 0$$

$$\nabla \times \mathbf{B} = \frac{4\pi}{c} \mathbf{J} + \frac{1}{c} \frac{\partial \mathbf{E}}{\partial t}$$

$$\nabla \cdot \mathbf{E} = 4\pi\rho$$

$$\nabla \times \mathbf{E} = -\frac{1}{c} \frac{\partial \mathbf{B}}{\partial t}$$

Vlasov-Maxwell Theory

Vlasov

$$\frac{\partial f_s}{\partial t} + \mathbf{v} \cdot \frac{\partial f_s}{\partial \mathbf{x}} + \frac{q_s}{m_s} \left(\mathbf{E} + \frac{\mathbf{v} \times \mathbf{B}}{c} \right) \cdot \frac{\partial f_s}{\partial \mathbf{v}} = 0$$

Maxwell

$$\begin{aligned}\nabla \cdot \mathbf{B} &= 0 & \nabla \times \mathbf{B} &= \frac{4\pi}{c} \mathbf{J} + \frac{1}{c} \frac{\partial \mathbf{E}}{\partial t} \\ \nabla \cdot \mathbf{E} &= 4\pi\rho & \nabla \times \mathbf{E} &= -\frac{1}{c} \frac{\partial \mathbf{B}}{\partial t}\end{aligned}$$

- **Coupled by first 2 moments** $\rho = \sum_s q_s \int f_s d\mathbf{v}$ $\mathbf{J} = \sum_s q_s \int \mathbf{v} f_s d\mathbf{v}$
- **Complete description of collisionless plasma**
- **Very difficult to solve - 6D phase space!**
- **Fluid description is much easier**

Fluid Description of Plasma

Density → $n_s = \int f_s d\mathbf{v}$

Velocity → $\mathbf{U}_s = \int \mathbf{v} f_s d\mathbf{v}$

Pressure → $\mathbf{P}_s = m_s \int (\mathbf{v} - \mathbf{U}_s)(\mathbf{v} - \mathbf{U}_s) f_s d\mathbf{v}$

Take velocity space moments of the Vlasov Equation:

Mass conservation → $\frac{\partial n_s}{\partial t} + \nabla \cdot (n_s \mathbf{U}_s) = 0$

Momentum conservation → $m_s n_s \frac{d\mathbf{U}_s}{dt} = -\nabla \cdot \mathbf{P}_s + q_s n_s \left(\mathbf{E} + \frac{\mathbf{U}_s \times \mathbf{B}}{c} \right)$

Closure Problem - Each equation contains higher order moment!

MHD Model

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{U}) = 0 \quad \rightarrow \quad \text{Mass conservation}$$

$$\rho \frac{d\mathbf{U}}{dt} = -\nabla P + \frac{\mathbf{J} \times \mathbf{B}}{c} \quad \rightarrow \quad \text{Momentum conservation}$$

$$\mathbf{E} + \frac{\mathbf{U} \times \mathbf{B}}{c} = \eta \mathbf{J} \quad \rightarrow \quad \text{Ohm's Law}$$

$$\frac{d}{dt} \left(\frac{P}{\rho^\gamma} \right) = 0 \quad \rightarrow \quad \text{Adiabatic equation of state}$$

$$\nabla \times \mathbf{B} = \frac{4\pi}{c} \mathbf{J} \quad \rightarrow \quad \text{Ampere's Law}$$

$$\nabla \times \mathbf{E} = -\frac{1}{c} \frac{\partial \mathbf{B}}{\partial t} \quad \rightarrow \quad \text{Faraday's Law}$$

$$\nabla \cdot \mathbf{B} = 0$$

$$\text{Ideal MHD} \rightarrow \eta = 0 \rightarrow \mathbf{U}_\perp = c \frac{\mathbf{E} \times \mathbf{B}}{B^2}$$

Evolution Equation for B Field

$$\frac{\partial \mathbf{B}}{\partial t} = \underbrace{\nabla \times (\mathbf{U} \times \mathbf{B})}_{\text{Convection}} + \underbrace{\frac{\eta c^2}{4\pi} \nabla^2 \mathbf{B}}_{\text{Diffusion}}$$

Magnetic
Diffusion
Coefficient

$$D_M = \frac{\eta c^2}{4\pi} \sim \frac{(\text{Length})^2}{\text{time}}$$

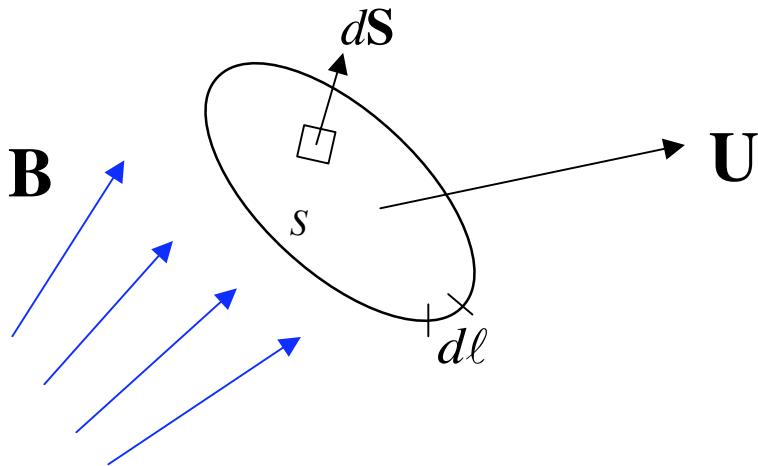
Time Scale for
Resistive Diffusion

$$\tau_R = \frac{L^2}{D_M}$$



$\tau_R \sim 10,000$ years!

Frozen-in Condition



Magnetic Flux

$$\psi = \int_s \mathbf{B} \bullet d\mathbf{S}$$

$$\frac{d\psi}{dt} = \int_s \frac{\partial \mathbf{B}}{\partial t} \bullet d\mathbf{S} - \oint (\mathbf{U} \times \mathbf{B}) \bullet d\ell$$

$$\frac{d\psi}{dt} = -c \oint \left(\mathbf{E} + \frac{\mathbf{U} \times \mathbf{B}}{c} \right) \bullet d\ell$$

Ohm's Law

$$\mathbf{E} + \frac{\mathbf{U} \times \mathbf{B}}{c} = \eta \mathbf{J}$$

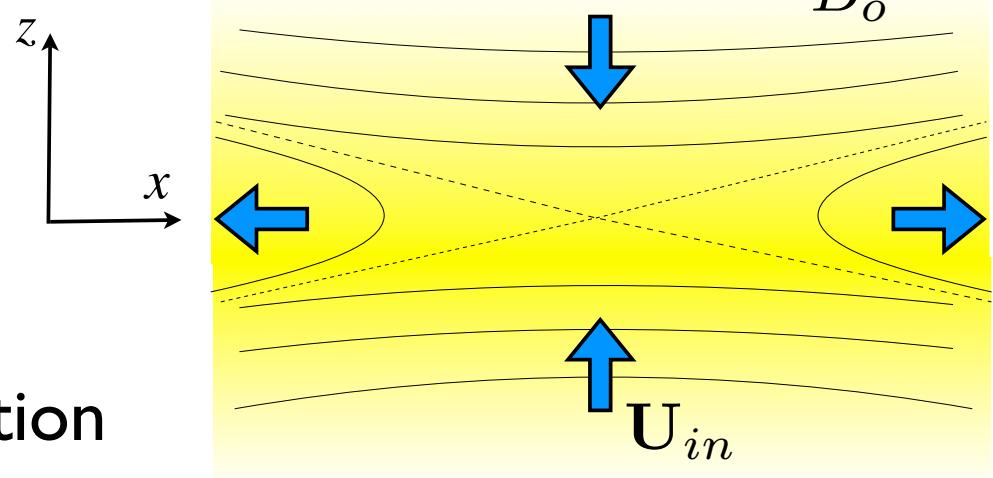
Ideal MHD $\rightarrow \frac{d\psi}{dt} = 0$

1. Magnetic field is “frozen” into plasma
2. No Topological changes

Define Basic Terms

- Reconnection rate

$$E_y = -\frac{1}{c} \frac{\partial A_y}{\partial t}$$



- Steady-state reconnection

$$\frac{\partial \mathbf{B}}{\partial t} = \frac{\partial \mathbf{E}}{\partial t} = 0 \quad \longrightarrow \quad E_y \text{ is spatially uniform in 2D}$$

- Collisionless reconnection

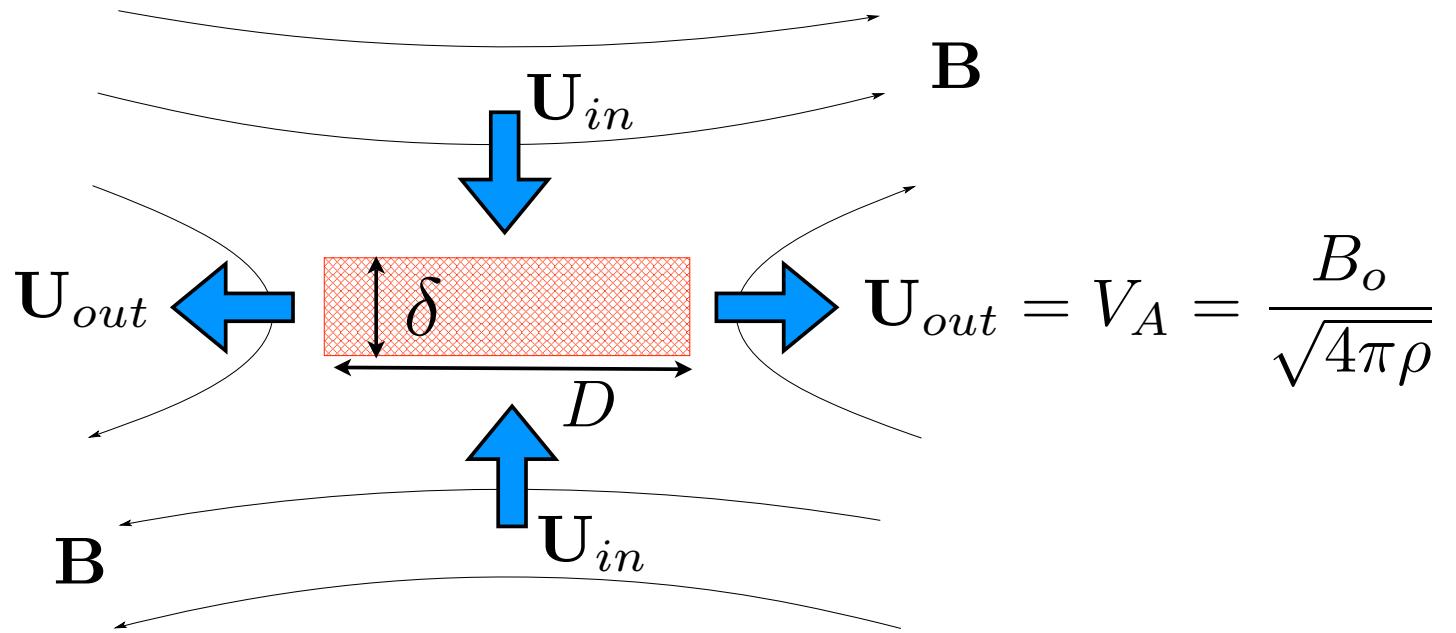
Dimensionless Reconnection Rate

$$E_R = \frac{\mathbf{U}_{in}}{V_A} = \frac{c E_y}{B_o V_A}$$

- Fast reconnection

Sweet and Parker's Model of Reconnection

1957 - 1958



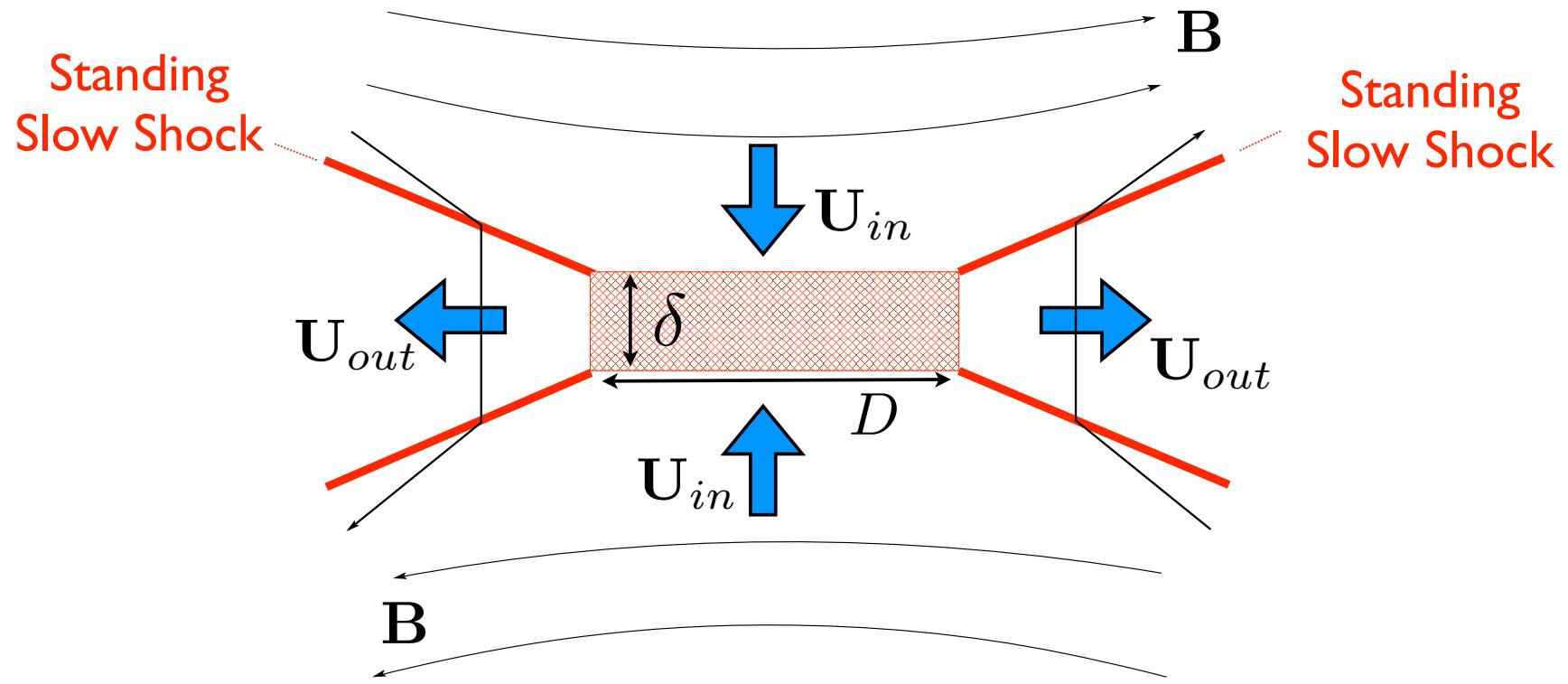
Reconnection
Rate

$$\frac{U_{in}}{V_A} = \frac{\delta}{D} = \frac{1}{\sqrt{S}}$$

$$S \equiv \frac{4\pi V_A D}{\eta c^2} = \frac{\tau_R}{\tau_A}$$

Lundquist Number

Petschek's Model - 1964

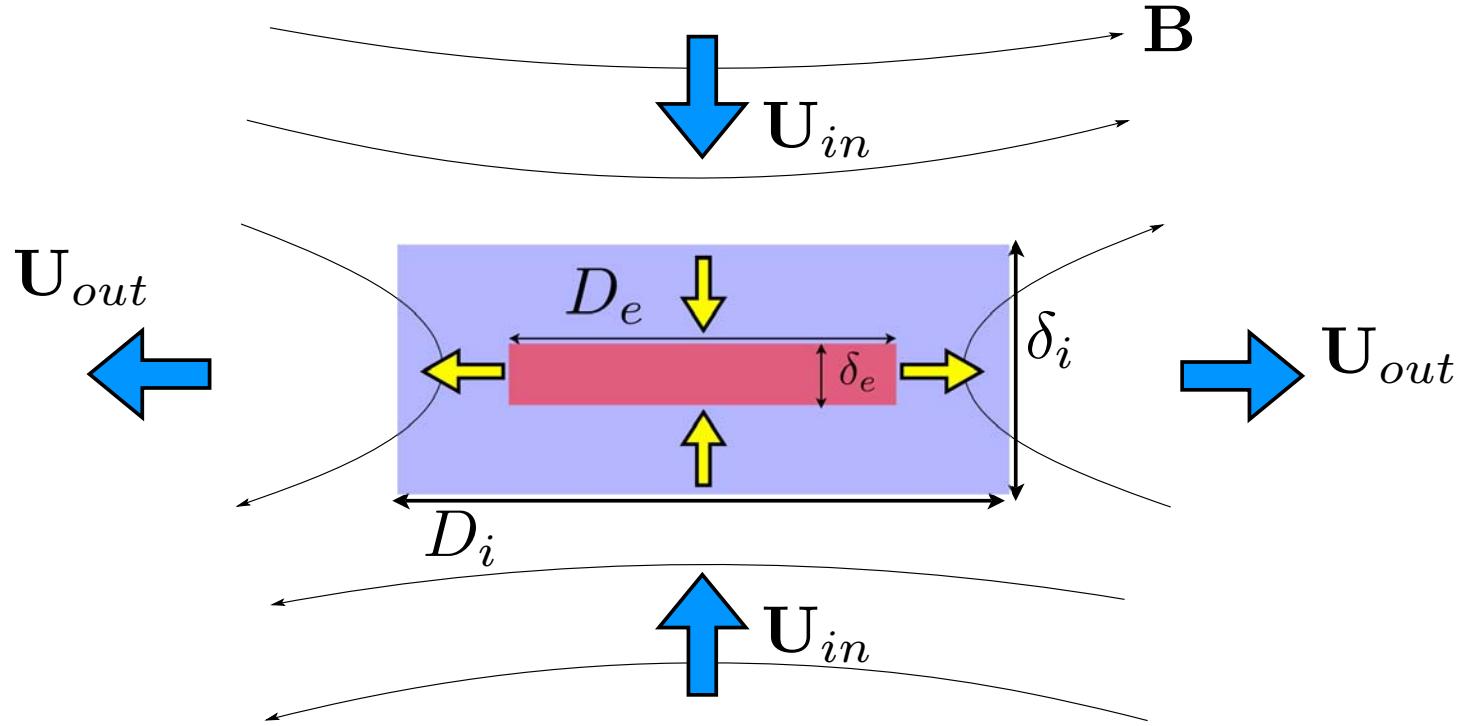


Reconnection
Rate

$$\frac{U_{in}}{V_A} \sim \frac{1}{\log(S)}$$

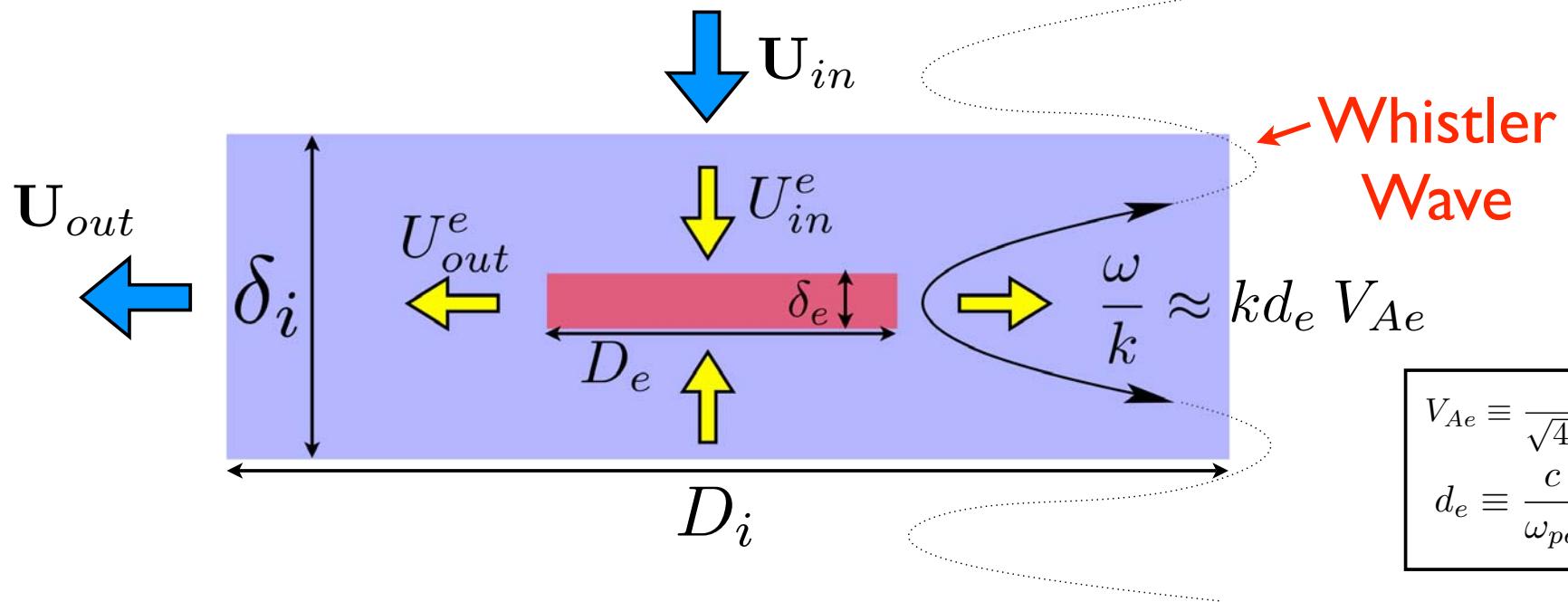
Consistent with observations
Widely accepted for 20+ years!

Hall Mediated Fast Reconnection 1995-2001



- $E + \frac{1}{c} U_i \times B \neq 0$ **Frozen-in condition violated for ions**
- $E + \frac{1}{c} U_e \times B \neq 0$ **Frozen-in condition violated for electrons**

Hall Mediated Reconnection



Electron Diffusion Region

$$D_e U_{in}^e = \delta_e U_{out}^e$$

$$U_{out}^e \sim \frac{d_e V_{Ae}}{\delta_e}$$

$$U_{in}^e \sim \frac{\delta_e}{D_e} \frac{d_e V_{Ae}}{\cancel{\delta_e}} \sim \frac{d_e V_{Ae}}{D_e}$$

$\cancel{\delta_e}$
cancel

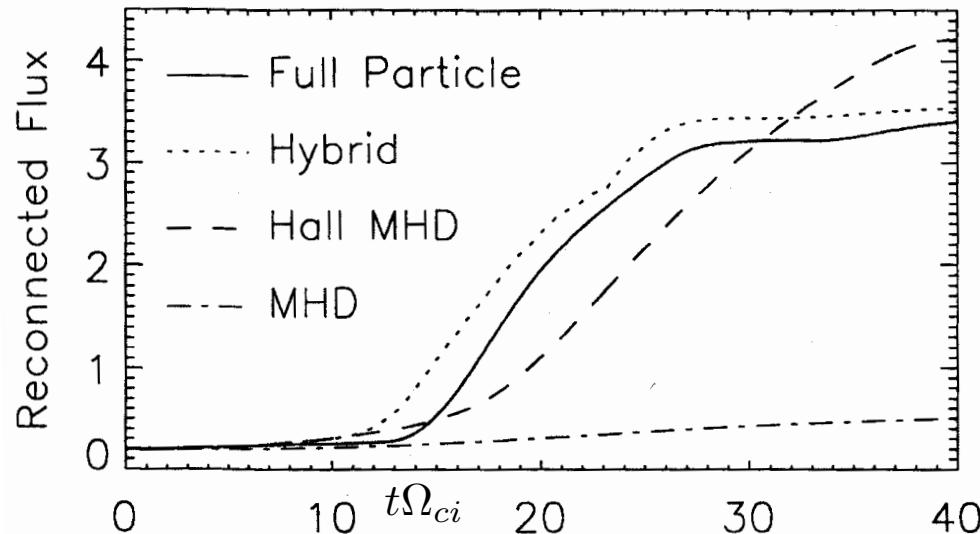
Ion Diffusion Region

$$D_i U_{in} = \delta_i U_{out}$$

$$U_{out} \leq V_{Ai}$$

$$\frac{U_{in}}{V_{Ai}} = \frac{\delta_i}{D_i}$$

GEM Challenge



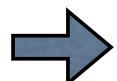
Birn, et al. 2001
Shay, et al. 2001
Pritchett, 2001
Hesse et al. 2001
Ma & Bhattacharjee, 2001
Otto, 2001

“Universal” Fast
Reconnection Rate

$$\frac{U_{in}}{V_A} \sim 0.1$$

Shay & Drake 1998
Rogers et al 2001
Shay et al 1999
Shay et al 2004
Huba & Rudakov 2004

Claim result is
independent of



1. Electron physics
2. Most plasma parameters
3. System size

Kinetic Simulation

Maxwell

$$\nabla \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t}$$

$$\nabla \cdot \vec{B} = 0$$

$$\nabla \times \vec{B} = \frac{4\pi}{c} \vec{J} + \frac{1}{c} \frac{\partial \vec{E}}{\partial t}$$

$$\nabla \cdot \vec{E} = 4\pi\rho$$

Vlasov

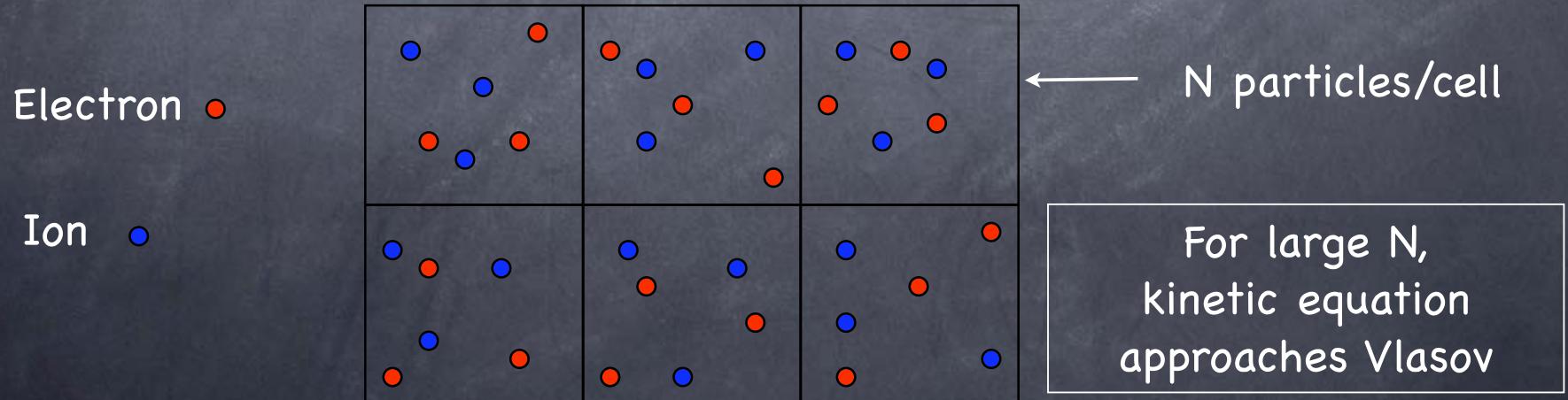
$$\frac{\partial f_s}{\partial t} + \vec{v} \cdot \frac{\partial f_s}{\partial \vec{x}} + \frac{q_s}{m_s} \left(\vec{E} + \frac{\vec{v} \times \vec{B}}{c} \right) \cdot \frac{\partial f_s}{\partial \vec{v}} = 0$$

- ⦿ Coupled by first 2 moments
- ⦿ Complete description of collisionless plasma
- ⦿ Possible to add collisions (difficult to do rigorously)
- ⦿ Easy to solve Maxwell's equations \vec{A} and ϕ
- ⦿ Vlasov is more difficult - Vlasov code vs PIC Code

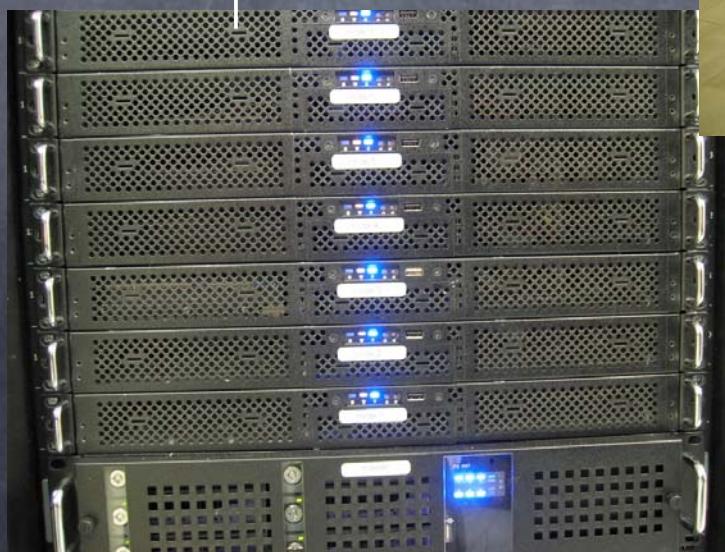
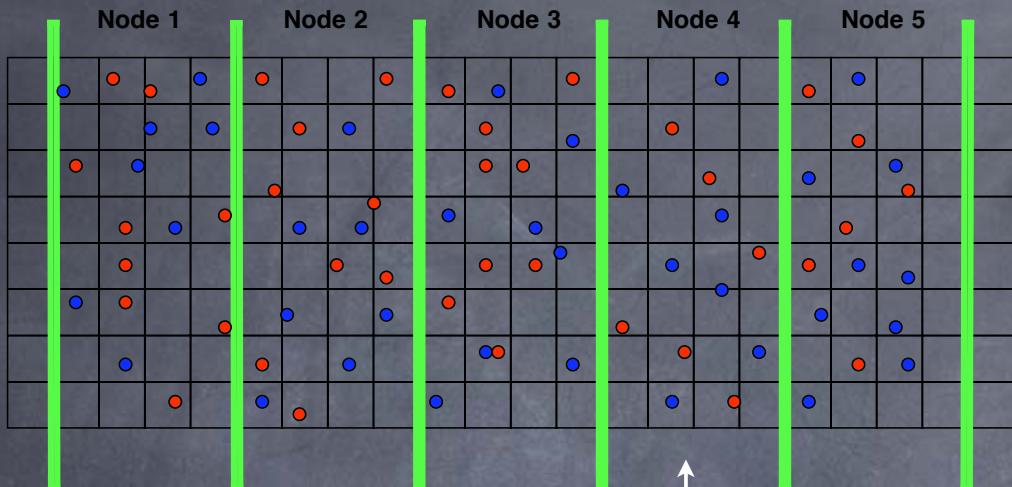
$$\rho = \sum_s q_s \int f_s d\vec{v} \quad \vec{J} = \sum_s q_s \int \vec{v} f_s d\vec{v}$$

PIC = Particle-in-cell

- Introduce “super-particles” - Lagrangian tracers
- Create spatial grid (cells)
- Interpolate position and velocity of particle onto grid $\longrightarrow \rho \vec{J}$
- Compute resulting E and B fields
- Push particles using these self-consistent fields
- Evolution of this system obeys a kinetic equation



Parallel Kinetic Simulation



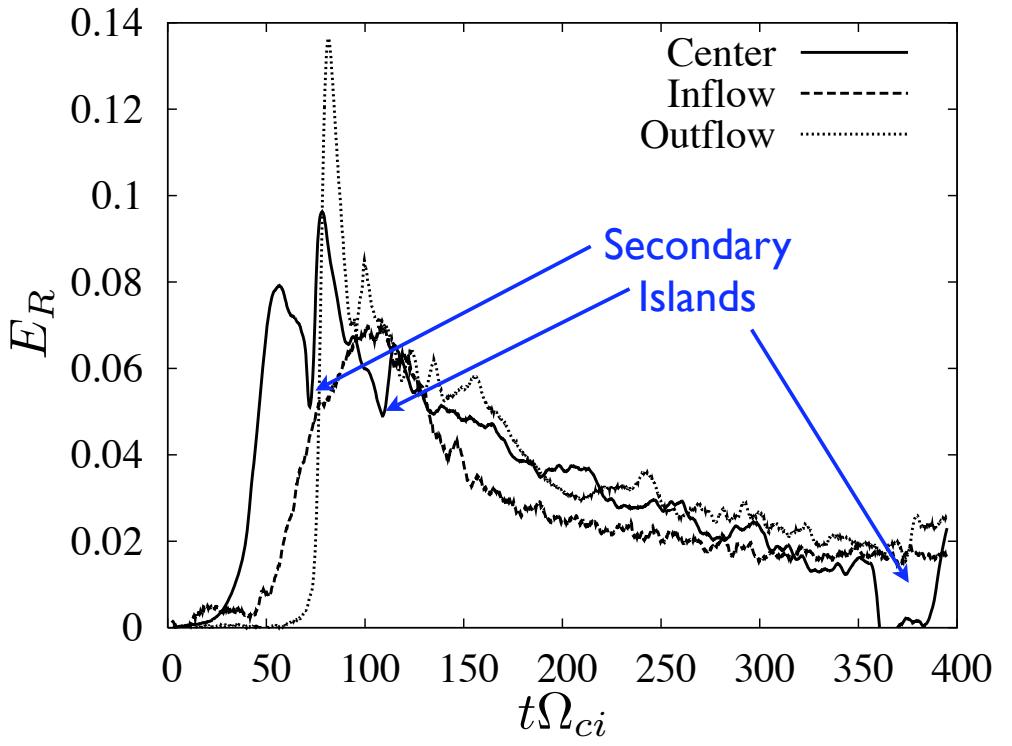
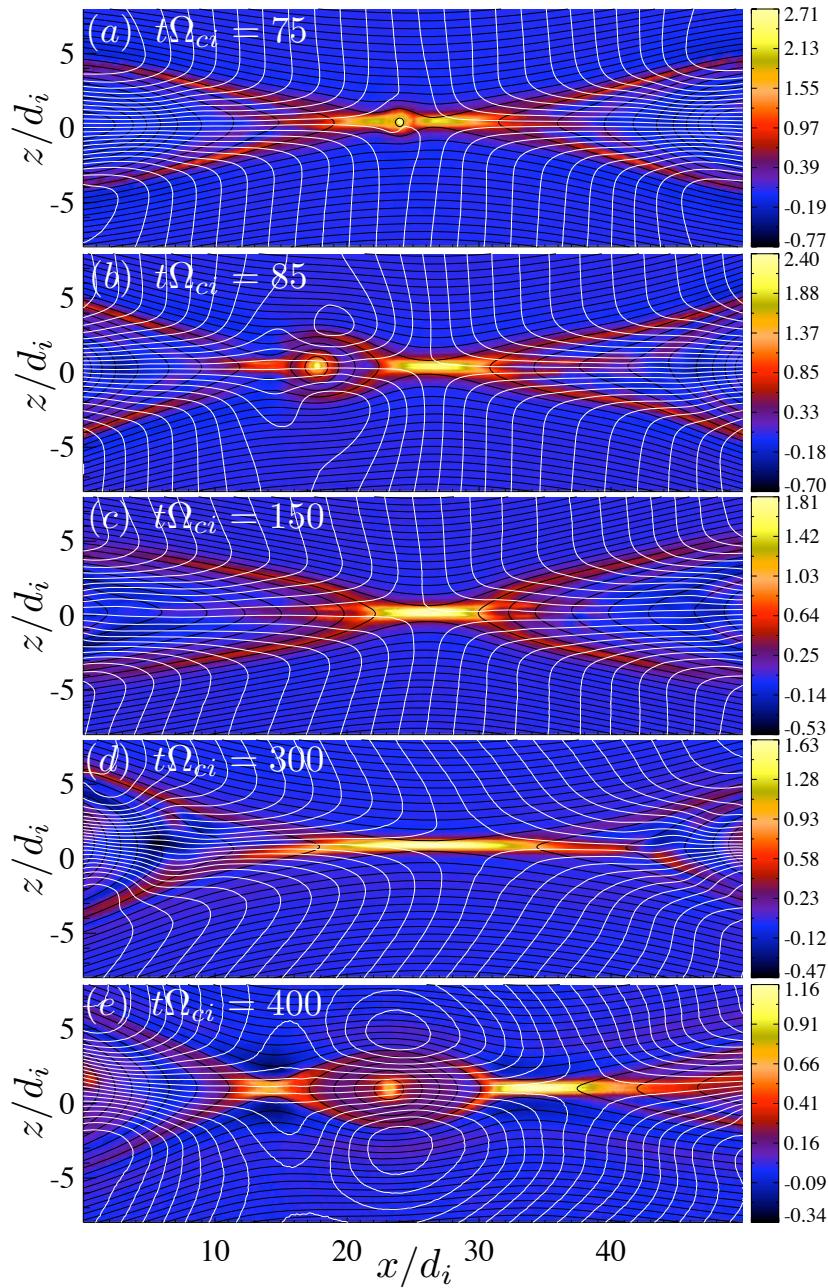
33 - SMP Nodes
2 Opteron CPU's in each
5 GB RAM - 165 GB Total
4 TB Raid Storage
Gigabit Ethernet

64 processor
Linux Cluster

Pitfalls of PIC Simulation

- Artificial simulation parameters → $\frac{m_i}{m_e} \quad \frac{\omega_{pe}}{\Omega_{ce}}$
- Simulations are mostly 2D
- Small system size
- Short simulation time
- **Periodic boundary conditions**

Large Open Boundary Case



Maximum length of electron diffusion region is limited by secondary island production

Future Work

- Essential to understand how these results scale with mass ratio and plasma parameters
- Evolution in 3D may be dramatically different due to large variety of possible plasma instabilities
- Unlikely to solve the problem with “brute force” computing alone - need theory, space observations, laboratory experiments
- To fully understand the physics of collisionless reconnection remains an enormous challenge